

Approximation theory for FEM

The question to be answered here is how big the error $u - u_h$ is ?

Here u_h belongs to given finite element space.

First, let's consider V_h = piecewise linear functions in R^2

For $k \in J_h$,

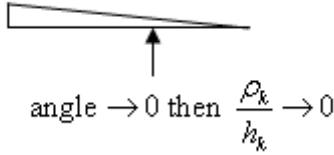
$$\begin{cases} h_k = \text{the diameter of } k = \text{longest side of } k \\ \rho_k = \text{the diameter of the circle inscribed in } k \\ h = \max_{h_k \in J_k} h_k \end{cases}$$

we assume \exists a constant β such that

$$\frac{\rho_k}{h_k} \geq \beta \quad \forall k \in J_h \quad (\text{meaning no triangle in } J_h \text{ is degenerated!})$$

Let N_i , $i = 1 \sim n$, be the nodes of J_h and define the interpolant

$$(\pi_h v)(N_i) = v(N_i) \quad \left(\begin{array}{l} \text{meaning } V_h \text{ consists of nodal interpolation} \\ \text{functions. This is not the case when using sin,} \\ \text{cos as basis functions.} \end{array} \right)$$



Theorem 1.

Let $k \in J_h$ be an element with vertices a^i , $i = 1 \sim 3$. Given $v \in C^0(k)$,

let the interpolant $\pi_h v \in P_1(k)$ be defined by $\pi_h v(a^i) = v(a^i)$, $i = 1 \sim 3$.

$$\text{Then (1)} \|v - \pi_h v\|_{L^\infty(k)} \leq 2h_k^2 \max_{|\hat{o}|=2} \|D^{\hat{o}} v\|_{L^\infty(k)}$$

Recall Theorem 7(*)

$$(2) \max_{|\hat{o}|=1} \|D^{\hat{o}}(v - \pi_h v)\|_{L^\infty(k)} \leq 6 \cdot \frac{h_k^2}{\rho_k} \max_{|\hat{o}|=2} \|D^{\hat{o}} v\|_{L^\infty(k)}$$

where $\|v\|_{L^\infty(k)} = \max_{x \in k} |v(x)|$

(*) The Sobolev inequality (2) C^0 can be embedded in H^2 in R^2 .

Now, since $|a_j^i - x_j| \leq h_k$, for $i = 1, 2, 3$, $j = 1, 2$,

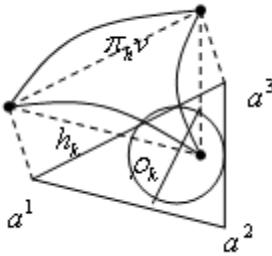
$$|R_i(x)| \leq 2h_k^2 \max_{|\partial|=2} \|D^\partial v\|_{L^\infty(k)}$$

Combine (1) & (2)

$$\begin{aligned} \Rightarrow \pi_h v &= \sum_{i=1}^2 v(x) \phi_i(x) + \sum_{i=1}^3 P_i(x) \phi_i(x) + \sum_{i=1}^3 R_i(x) \phi_i(x) \\ &= v(x) \sum_{i=1}^3 \phi_i(x) + \sum_{i=1}^3 P_i(x) \phi_i(x) + \sum_{i=1}^3 R_i(x) \phi_i(x) \quad -(3) \end{aligned}$$

Observation (1) $\sum_{i=1}^3 \phi_i(x) = 1$

$$\Rightarrow |\pi_h v - v(x)| \leq \left| \sum_{i=1}^3 P_i(x) \phi_i(x) \right| + \left| \sum_{i=1}^3 R_i(x) \phi_i(x) \right| \quad -(*)$$



proof:

Let ϕ_i , $i = 1 \sim 3$, be the basis functions

$$\text{For } w \in P_1(k), \text{ we have } w(x) = \sum_{i=1}^3 W(a^i) \phi_i(x) \quad \text{for } x \in k$$

$$\Rightarrow (\pi_h v)(x) = \sum_{i=1}^3 v(a^i) \phi_i(x) \quad \text{for } x \in k \quad -(1)$$

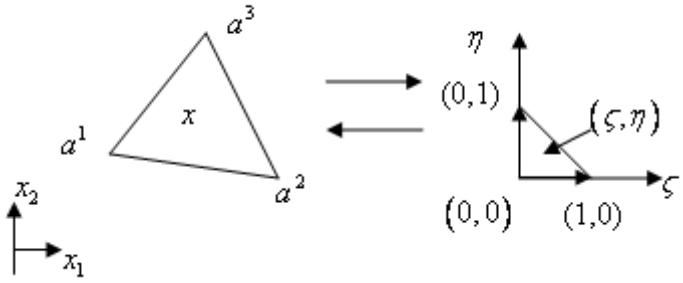
By Taylor expansion, one has $x = (x_1, x_2)$, $y = (y_1, y_2)$

$$v(y) = v(x) + \sum_{j=1}^2 \frac{\partial v(x)}{\partial x_j} (y_j - x_j) + R(x, y)$$

$$\text{here } R(x, y) = \frac{1}{2} \sum_{i,j=1}^2 \frac{\partial^2 v(\xi)}{\partial x_i \partial x_j} (y_i - x_i)(y_j - x_j)$$

Now choose $y = a^i$, we have $v(a^i) = v(x) + P_i(x) + R_i(x) \quad -(2)$

$$\begin{cases} P_i(x) = \sum_{j=1}^2 \frac{\partial v}{\partial x_j} (a_j^i - x_j) \quad (a^i = (a_1^i, a_2^i)) \\ \quad \underbrace{_{= \nabla v \cdot (a^i - x)}} \\ R_i(x) = R(x, a^i) \end{cases}$$



Now consider the parametrization

$$\text{It can be shown } x = (1 - \xi - \eta) a^1 + \xi a^2 + \eta a^3$$

$$\begin{aligned} \text{and } \sum_{i=1}^3 P_i(x) \phi_i(x) &= \sum_{i=1}^3 \nabla v((a^i - x) \cdot \phi_i(x)) \\ &= \sum_{\text{parametrization}} \nabla v \left(a^i - [(1 - \xi - \eta) a^1 + \xi a^2 + \eta a^3] \cdot \phi_i(\xi, \eta) \right) \\ &= \nabla v \cdot \left\{ \begin{aligned} &\left(a^1 - [(1 - \xi - \eta) a^1 + \xi a^2 + \eta a^3] \right) \cdot (1 - \xi - \eta) + \\ &\left(a^2 - [(1 - \xi - \eta) a^1 + \xi a^2 + \eta a^3] \right) \cdot \xi + \\ &\left(a^3 - [(1 - \xi - \eta) a^1 + \xi a^2 + \eta a^3] \right) \cdot \eta \end{aligned} \right\} \\ &= 0 \quad (\text{By direct computation}) \quad -(4) \end{aligned}$$

Therefore,

$$\begin{aligned} (*) \Rightarrow |\pi_h v - v| &\leq \left| \sum_{i=1}^3 R_i(x) \phi_i(x) \right| \leq \sum_{i=1}^3 \max_{x \in k} |R_i(x)| |\phi_i(x)| \\ &\leq \max_{\substack{x \in k \\ i=\{2\}}} |R_i(x)| \leq 2h_k^2 \max_{|\hat{\theta}|=2} \|D^\hat{\theta} v\|_{L^\infty(k)} \end{aligned}$$

(1) is proved.

To prove (2), let's differentiate (3)

$$\stackrel{(3)&(4)}{\Rightarrow} D^{(\partial)} \pi_h v = D^{(\partial)} v(k) \left(\sum_{i=1}^3 \phi_i(x) \right) + v(x) \sum_{i=1}^3 D^{(\partial)} \phi_i(x) + D^{(\partial)} \sum_{i=1}^3 R_i(x) \phi_i(x) \quad -(5)$$

it can be shown

$$\sum_{i=1}^3 D^{(\partial)} \phi_i(x) = 0 \quad \xleftrightarrow{\text{by parametrization}} \quad \begin{cases} \frac{\partial}{\partial \zeta} (1 - \zeta - \eta) + \frac{\partial}{\partial \zeta} (\zeta) + \frac{\partial}{\partial \zeta} (\eta) = 0 \\ \frac{\partial}{\partial \eta} (1 - \zeta - \eta) + \frac{\partial}{\partial \eta} (\zeta) + \frac{\partial}{\partial \eta} (\eta) = 0 \end{cases}$$

This implies

$$D^{(\partial)} \pi_h v - D^{(\partial)} v = D^{(\partial)} \sum_{i=1}^3 R_i(x) \phi_i(x) = D^{(\partial)} \sum_{i=1}^3 (a^i - x)^T [H] (a^i - x) \cdot \phi_i(x)$$

$$\text{here } [H] = \left[\frac{\partial^2 v(\xi)}{\partial x_i \partial x_j} \right].$$

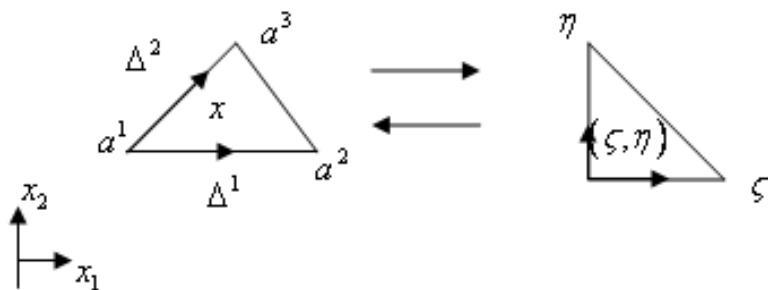
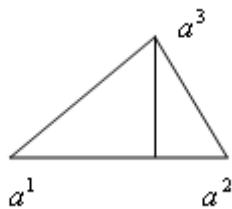
$$\because D^{(\partial)} ((a^i - x) H (a^i - x)) = 2[H] (a^i - x)$$

$$\Rightarrow \sum_{i=1}^3 (D^{(\partial)} R_i(x)) \phi_i(x) = 2[H] \sum_{i=1}^3 (a^i - x) \cdot \phi_i(x) = 0 \quad (\text{by (4)})$$

$$\Rightarrow |D^{(\partial)} \pi_h v - D^{(\partial)} v| = \left| \sum_{i=1}^3 R_i(x) D^{(\partial)} \phi_i(x) \right| \leq \sum_{i=1}^3 |R_i(x)| |D^{(\partial)} \phi_i(x)| \quad -(7)$$

Again, by parametrization, we can show

$$|D^{(\partial)} \phi_i(x)| = \begin{cases} \frac{\partial \zeta}{\partial x_1} \cdot \frac{\partial \phi_i}{\partial \zeta} + \frac{\partial \eta}{\partial x_1} \cdot \frac{\partial \phi_i}{\partial \eta} \\ \frac{\partial \zeta}{\partial x_2} \cdot \frac{\partial \phi_i}{\partial \zeta} + \frac{\partial \eta}{\partial x_2} \cdot \frac{\partial \phi_i}{\partial \eta} \end{cases} = \begin{bmatrix} \frac{\partial \zeta}{\partial x_1} & \frac{\partial \eta}{\partial x_1} \\ \frac{\partial \zeta}{\partial x_2} & \frac{\partial \eta}{\partial x_2} \end{bmatrix} \begin{bmatrix} \frac{\partial \phi_i}{\partial \zeta} \\ \frac{\partial \phi_i}{\partial \eta} \end{bmatrix}$$



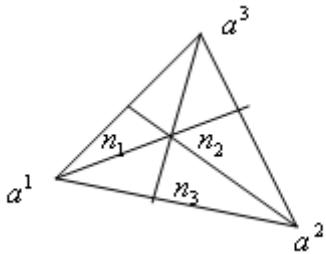
Recall:

$$T(\zeta, \eta) = \left(\underbrace{a^2 - a^1}_{\Delta^1}, \underbrace{a^3 - a^1}_{\Delta^2} \right) \begin{pmatrix} \zeta \\ \eta \end{pmatrix} + \vec{a}_1 = \vec{x}$$

$$\Rightarrow \begin{pmatrix} \zeta \\ \eta \end{pmatrix} = \frac{1}{|\Delta|} \begin{pmatrix} \Delta_2^2 & -\Delta_1^2 \\ -\Delta_2^1 & \Delta_1^1 \end{pmatrix} (\vec{x} - \vec{a}^1)$$

(Jacobian)

$$\Rightarrow \begin{pmatrix} \frac{\partial \zeta}{\partial x_1} & \frac{\partial \zeta}{\partial x_2} \\ \frac{\partial \eta}{\partial x_1} & \frac{\partial \eta}{\partial x_2} \end{pmatrix} = \frac{1}{|\Delta|} \begin{pmatrix} \Delta_2^2 & -\Delta_2^1 \\ -\Delta_1^2 & \Delta_1^1 \end{pmatrix}$$



$$\left\| \frac{1}{|\Delta|} \begin{pmatrix} \Delta_2^2 & -\Delta_2^1 \\ -\Delta_1^2 & \Delta_1^1 \end{pmatrix} \begin{pmatrix} \frac{\partial \phi_2}{\partial \zeta} \\ \frac{\partial \phi_2}{\partial \eta} \end{pmatrix} \right\| = \frac{1}{|\Delta|} \left\| \begin{pmatrix} \Delta_2^2 \\ -\Delta_1^2 \end{pmatrix} \right\| = \frac{\|a^3 - a_1\|}{|\Delta|} = \frac{1}{h_2}$$

$$\left\| \frac{1}{|\Delta|} \begin{pmatrix} \Delta_2^2 & -\Delta_2^1 \\ -\Delta_1^2 & \Delta_1^1 \end{pmatrix} \begin{pmatrix} \frac{\partial \phi_3}{\partial \zeta} \\ \frac{\partial \phi_3}{\partial \eta} \end{pmatrix} \right\| = \frac{1}{|\Delta|} \left\| \begin{pmatrix} -\Delta_2^1 \\ \Delta_1^1 \end{pmatrix} \right\| = \frac{\|a^2 - a_1\|}{|\Delta|} = \frac{1}{h_3}$$

$$\begin{aligned} \left\| \frac{1}{|\Delta|} \begin{pmatrix} \Delta_2^2 & -\Delta_2^1 \\ -\Delta_1^2 & \Delta_1^1 \end{pmatrix} \begin{pmatrix} \frac{\partial \phi_1}{\partial \zeta} \\ \frac{\partial \phi_1}{\partial \eta} \\ \begin{pmatrix} 1 \\ -1 \\ -1 \end{pmatrix} \end{pmatrix} \right\| &= \frac{1}{|\Delta|} \left\| \begin{pmatrix} \Delta_2^2 - \Delta_2^1 \\ \Delta_1^1 - \Delta_1^2 \\ \Delta_1^1 - \Delta_2^1 \end{pmatrix} \right\| = \frac{1}{|\Delta|} = \frac{1}{h_2} \left\| \begin{pmatrix} \Delta_2^1 - \Delta_2^2 \\ \Delta_1^1 - \Delta_1^2 \end{pmatrix} \right\| \\ &= \frac{\|a^2 - a^3\|}{|\Delta|} = \frac{1}{h_1} \end{aligned}$$

$$\Rightarrow \|D^{(\partial)} \phi_i\|_k < \frac{1}{\rho_k}$$

As a result, (7) becomes

$$|D^{(\partial)}\pi_h v - D^{(\partial)}v| \leq \frac{1}{\rho_k} \sum_{i=1}^3 |R_i(x)| \leq 6 \frac{h_k^2}{\rho_k} \max_{|\partial|=2} \|D^\partial v\|_{L^\infty(k)}$$

This proves (2).

Let's define the semi-norm

$$|v|_{H_2^r} = \left(\sum_{|\partial|=r} \int |D^\partial v|^2 \right)^{\frac{1}{2}}$$

(because $|v|_{H_2^r} = 0$ even if $v \neq 0$ (e.g. $v \equiv 1, r \geq 1$))

Poincare inequality ($\|u\|_{L^2} < \|\nabla u\|_{L^2}$) shows that it is suffice to estimate the H_2^r semi-norm for obtaining the upper bound of the H_2^r -norm.

Theorem 2. Under the assumption of theorem 1, there is a constant c such that

$$\|v - \pi_h v\|_{L_2(k)} \leq ch_k^2 |v|_{H^2(k)}$$

$$|v - \pi_h v|_{H^1(k)} \leq c \frac{h_k^2}{\rho_k} |v|_{H^2(k)}$$

(Pf: skip.)

Global error estimation

Theorem 3. Under the assumption of Theorem 1,

$$\text{we have } \|v - \pi_h v\|_{L^2(\Omega)} \leq ch^2 |u|_{H^2(\Omega)}$$

$$|v - \pi_h v|_{H^1(\Omega)} \leq ch |u|_{H^2(\Omega)}$$

$$\text{Pf: } \|v - \pi_h v\|_{L^2(\Omega)}^2 = \sum_{k \in J_n} \|v - \pi_h v\|_{L^2(\Omega)}^2 \leq \sum_{k \in J_n} c^2 h_k^4 |u|_{H^2(k)}^2 \leq c^2 h^4 |u|_{H^2(\Omega)}^2$$

$$|v - \pi_h v|_{H^1}^2 \leq \sum_{k \in J_n} c^2 \frac{h_k^4}{\rho_k^2} |u|_{H^2(k)}^2 \stackrel{\left(\frac{\rho_k}{h_k} \geq \beta\right)}{\leq} \sum_{k \in J_n} c^2 \frac{h_k^2}{\beta^2} |u|_{H^2(k)}^2 \leq c(\beta) h^2 |u|_{H^2(\Omega)}^2$$

Remark (2) error of FEM solution satisfies

$$\begin{aligned} \underbrace{\|u - u_h\|}_{\text{energy norm}} &\stackrel{\text{(Cea)}}{\leq} \min_{v \in V_h} \|u - v\| \stackrel{\text{(Poincare)}}{\leq} c \min_{v \in V_h} |u - v|_{H^1} \stackrel{\text{(interpolation)}}{\leq} \underset{\text{error}}{\frac{c}{\rho_k}} ch |u|_{H^2(\Omega)} \end{aligned}$$

Remark (3) Error estimation in L^2 -norm

Consider $a(u, v) = \int A \nabla u \cdot \nabla v + cuv dx$. Suppose u and u_h are the weak

solution and the FEM solution respectively. $\begin{cases} \text{i.e. } a(u, v) = F(v) & v \in V \\ a(u_h, v) = F(v) & v \in V_h \end{cases}$

Let $e = u - u_h$. Clearly one has $a(u - u_h, v) = 0 \Rightarrow a(e, v) = 0 \quad v \in V_h$

Consider φ the solution of $\begin{cases} a(v, \varphi) = (v, e) \\ \varphi|_{\Omega} = 0 \end{cases} \quad (*)$

$(*)$ is called the adjoint problem)

For any $\varphi_h \in V_h$ and $v = u - u_h$, we have

$$\begin{aligned} a(v, \varphi - \varphi_h) &= a(u - u_h, \varphi) + a(u - u_h, \varphi_h) \\ &= a(u - u_h, \varphi) = \langle u - u_h, e \rangle = a(u - u_h, \varphi - \varphi_h) \\ \Rightarrow \|u - u_h\|_{L^2}^2 &\stackrel{\text{Holdering}}{\leq} c \|u - u_h\|_{H^1(\Omega)} \|\varphi - \varphi_h\|_{H^1(\Omega)} \\ &\leq \tilde{c} h \|u - u_h\|_{H^1} |\varphi|_{H^2(\Omega)} \quad (\text{by Remark (2)}) \end{aligned}$$

When the adjoint problem $(*)$ has H^2 -regularity (for example Ω is convex)

we have $|\varphi|_{H^2(\Omega)} < \|e\|_{L^2(\Omega)}$

Using this inequality, one has

$$\begin{aligned} \|u - u_h\|_{L^2(\Omega)}^2 &\leq \tilde{c} h \|u - u_h\|_{H^1(\Omega)} \cdot \|u - u_h\|_{L^2(\Omega)} \\ \Rightarrow \|u - u_h\|_{L^2(\Omega)} &\leq \tilde{c} h^2 |u|_{H^2(\Omega)} \quad (\text{by remark (2)}) \end{aligned}$$

The above argument is the so-called "duality argument".

Remark (4): Interpolation with poly of higher degree $r \geq 1$

$$\begin{aligned} \|u - \pi_h u\|_{L^2(\Omega)} &\leq ch^{r+1} |u|_{H^{r+1}(\Omega)} \\ \|u - \pi_h u\|_{H^1(\Omega)} &\leq ch^r |u|_{H^{r+1}(\Omega)} \\ \|u - \pi_h u\|_{H^2(\Omega)} &\leq ch^{r-1} |u|_{H^{r+1}(\Omega)} \end{aligned}$$

If u doesn't have regularity of H^{r+1} , then suppose $u \in H^s$, $1 \leq s \leq r+1$,

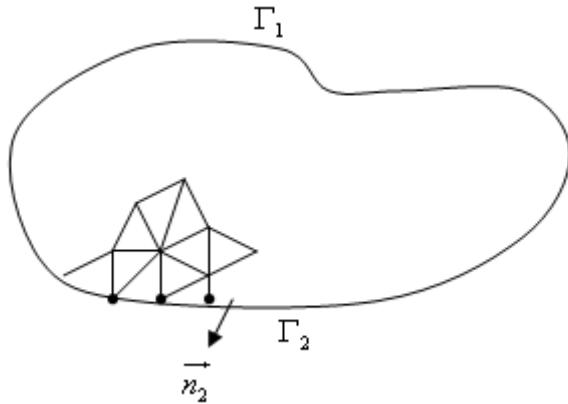
we have $\|u - \pi_h u\|_{L^2(\Omega)} \leq ch^s |u|_{H^s(\Omega)}$

$$\|u - \pi_h u\|_{H^1(\Omega)} \leq ch^{s-1} |u|_{H^s(\Omega)}$$

General FEM Procedure for solving P.D.E

Consider the following P.D.E

$$\begin{cases} -\operatorname{div}(A_{2 \times 2} \cdot \nabla u) + B \nabla u + cu = f \\ u|_{\Gamma_1} = g; \quad A \nabla u \cdot \vec{n} = r \end{cases}$$



Step 1: write down the weak formulation

$$\begin{aligned}
 & \int_{\Omega} v \left[(-\operatorname{div}(A \nabla u) + B \cdot \nabla u + cu) \right] dx = \int_{\Omega} v f dx, \text{ here } v|_{\Gamma_1} = 0 \\
 & \xrightarrow{\text{divergence theorem}} \underbrace{\int_{\Omega} \nabla v \cdot A \nabla u}_{(I)} + \underbrace{\int_{\Omega} v (B \cdot \nabla u)}_{(II)} + \underbrace{\int_{\Omega} c v u}_{(III)} \\
 & = \int_{\Omega} v \cdot f + \int_{\Gamma_2} v (A \nabla u) \cdot \overrightarrow{n}_2 - (8) \\
 & = \underbrace{\int_{\Omega} v \cdot f dx}_{(IV)} + \underbrace{\int_{\Gamma_2} v \cdot r ds}_{(V)}
 \end{aligned}$$

normal vector
of Γ_2

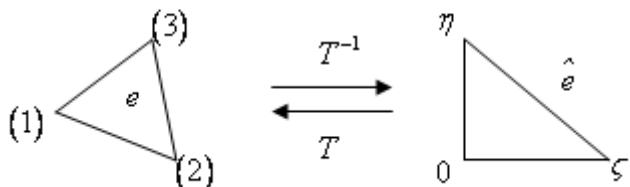
Step 2. FEM discretization

(a) generate mesh J_h for Ω

(b) consider (2)

$$\int_{\Omega} \nabla v \cdot A \nabla u dx = \sum_{e \in J_h(\Omega)} \int_e \nabla v \cdot A \nabla u dx$$

parametrization for e :



$$\begin{aligned}
\begin{pmatrix} x \\ y \end{pmatrix} &= \begin{pmatrix} x \\ y \end{pmatrix}^{(1)} + T(\varsigma, \eta) \\
&= \begin{pmatrix} x \\ y \end{pmatrix}^{(1)} (1 - \varsigma - \eta) + \begin{pmatrix} x \\ y \end{pmatrix}^{(2)} \varsigma + \begin{pmatrix} x \\ y \end{pmatrix}^{(3)} \eta \\
&= \begin{pmatrix} x \\ y \end{pmatrix}^{(1)} \underbrace{\begin{pmatrix} \Delta_2 x & \Delta_3 x \\ \Delta_2 y & \Delta_3 y \end{pmatrix}}_{=\Phi} \begin{pmatrix} \varsigma \\ \eta \end{pmatrix} + \begin{pmatrix} x \\ y \end{pmatrix}^{(1)} \\
\Rightarrow \begin{pmatrix} \varsigma \\ \eta \end{pmatrix} &= T^{-1} \begin{pmatrix} x \\ y \end{pmatrix} = \Phi^{-1} \begin{pmatrix} x \\ y \end{pmatrix} + \Phi^{-1} \begin{pmatrix} x \\ y \end{pmatrix}^{(1)}
\end{aligned}$$

We have

$$\begin{aligned}
\frac{\partial u}{\partial x} &= \frac{\partial u}{\partial \varsigma} \frac{\partial \varsigma}{\partial x} + \frac{\partial u}{\partial \eta} \frac{\partial \eta}{\partial x} \\
\frac{\partial u}{\partial y} &= \frac{\partial u}{\partial \varsigma} \frac{\partial \varsigma}{\partial y} + \frac{\partial u}{\partial \eta} \frac{\partial \eta}{\partial y} \\
\Rightarrow \begin{pmatrix} \frac{\partial u}{\partial x} \\ \frac{\partial u}{\partial y} \end{pmatrix} &= \begin{pmatrix} \frac{\partial \varsigma}{\partial x} & \frac{\partial \eta}{\partial x} \\ \frac{\partial \varsigma}{\partial y} & \frac{\partial \eta}{\partial y} \end{pmatrix} \begin{pmatrix} \frac{\partial u}{\partial \varsigma} \\ \frac{\partial u}{\partial \eta} \end{pmatrix} = (\Phi^{-1})^T \begin{pmatrix} \frac{\partial u}{\partial \varsigma} \\ \frac{\partial u}{\partial \eta} \end{pmatrix}
\end{aligned}$$

$$\text{Let } \Phi = (\Phi^{-1})^T$$

Choose FEM spaces and discretizatin

$$(c) \int_e \nabla v A_{2 \times 2} \nabla u dx = \int_e \left[\frac{\partial v}{\partial \varsigma} \quad \frac{\partial v}{\partial \eta} \right] \Phi^T A(\Phi) \begin{pmatrix} \frac{\partial u}{\partial \varsigma} \\ \frac{\partial u}{\partial \eta} \end{pmatrix} J d\varsigma d\eta$$

Let $\phi_i \sim \phi_m$ be the basis functions on the reference element satisfying $\phi_i(x_j) = \delta_{ij}$ and $v_i \sim v_m$ be the nodal values

$$\begin{aligned}
v &= \sum_{i=1}^m v_i \phi_i, \quad u = \sum_{j=1}^m u_j \phi_j \\
\Rightarrow \begin{pmatrix} \frac{\partial u}{\partial \varsigma} \\ \frac{\partial u}{\partial \eta} \end{pmatrix} &= \underbrace{\begin{bmatrix} \frac{\partial \phi_1}{\partial \varsigma} & \dots & \frac{\partial \phi_m}{\partial \varsigma} \\ \frac{\partial \phi_1}{\partial \eta} & \dots & \frac{\partial \phi_m}{\partial \eta} \end{bmatrix}_{2 \times m}}_D \begin{bmatrix} u_1 \\ \vdots \\ u_m \end{bmatrix} \\
\left[\frac{\partial v}{\partial \varsigma}, \frac{\partial v}{\partial \eta} \right] &= [v_1 \cdots v_m] D_{m \times 2}^T
\end{aligned}$$

$$\int_e \nabla v A \nabla u dx = \int_{\hat{e}} [v_1 \dots v_m] [D^T \cdot \Phi^T A(\Phi) D |J| \begin{bmatrix} u_1 \\ \vdots \\ u_m \end{bmatrix}] dx$$

$$= \vec{v}_e \cdot \vec{K}_e \cdot \vec{u}_e \quad -(9)$$

here $\vec{v}_e = [v_1 \dots v_m]$, $\vec{u}_e = [u_1 \dots u_m]$,

$$K_e = \int_{\hat{e}} D^T \Phi^T A(\Phi) D |J| d\zeta d\eta \quad \leftarrow \text{need Numerical Integration}$$

K_e : stiffness matrix.

Notes: Closer look at computing Jacobian $|J|$

$$|J| = \det \begin{pmatrix} \frac{\partial x}{\partial \zeta} & \frac{\partial y}{\partial \zeta} \\ \frac{\partial x}{\partial \eta} & \frac{\partial y}{\partial \eta} \end{pmatrix}$$

let $\varphi_1 = 1 - \zeta - \eta$, $\varphi_2 = \zeta$, $\varphi_3 = \eta$ (basis of the mapping functions)

one has $x = x_1 \varphi_1 + x_2 \varphi_2 + x_3 \varphi_3$

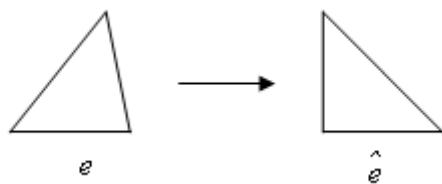
$$y = y_1 \varphi_1 + y_2 \varphi_2 + y_3 \varphi_3$$

$$\Rightarrow \frac{\partial x}{\partial \zeta} = \sum_{i=1}^3 x_i \frac{\partial \varphi_i}{\partial \zeta}, \quad \frac{\partial y}{\partial \zeta} = \sum_{i=1}^3 \frac{\partial \varphi_i}{\partial \zeta} y_i$$

$$\frac{\partial x}{\partial \eta} = \sum_{i=1}^3 x_i \frac{\partial \varphi_i}{\partial \eta}, \quad \frac{\partial y}{\partial \eta} = \sum_{i=1}^3 \frac{\partial \varphi_i}{\partial \eta} y_i$$

$$\Rightarrow |J| = \begin{vmatrix} \vec{\varphi} \circ \vec{x} & \vec{\varphi} \circ \vec{y} \\ \frac{\partial \vec{\varphi}}{\partial \zeta} \circ \vec{x} & \frac{\partial \vec{\varphi}}{\partial \zeta} \circ \vec{y} \\ \frac{\partial \vec{\varphi}}{\partial \eta} \circ \vec{x} & \frac{\partial \vec{\varphi}}{\partial \eta} \circ \vec{y} \end{vmatrix} \quad \begin{aligned} \vec{\varphi} &= \langle \varphi_1, \varphi_2, \varphi_3 \rangle \\ \vec{x} &= \langle x_1, x_2, x_3 \rangle \\ \vec{y} &= \langle y_1, y_2, y_3 \rangle \end{aligned}$$

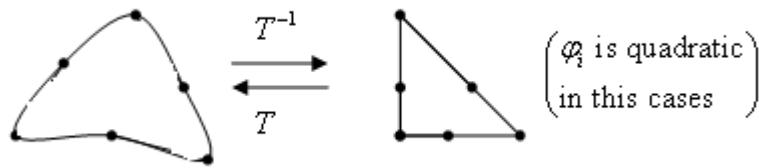
\circ : inner product.



In linear basis function, $|J|$ is a constant.

More general case, let $\varphi_1 \sim \varphi_m$ be basis of the mapping function

$$\begin{aligned} \begin{pmatrix} x \\ y \end{pmatrix} &= \sum_{i=1}^m \begin{pmatrix} x_i \\ y_i \end{pmatrix} \varphi_i \\ \Rightarrow |J| &= \begin{vmatrix} \frac{\partial \vec{\varphi}}{\partial \zeta} \circ \vec{x} & \frac{\partial \vec{\varphi}}{\partial \zeta} \circ \vec{y} \\ \frac{\partial \vec{\varphi}}{\partial \eta} \circ \vec{x} & \frac{\partial \vec{\varphi}}{\partial \eta} \circ \vec{y} \end{vmatrix} \rightarrow \text{Jacobian needs to be approximated in numerical integration.} \end{aligned}$$



(d) Similar discretization process applied on

$$\begin{aligned} (II) \int_{\Omega} v(B \cdot \nabla u) d\Omega &= \sum_{e \in J_h} \int_e v \underbrace{B}_{\substack{\uparrow \\ (\text{vector})}} \cdot \nabla u d\eta \\ \int_e v B \cdot \nabla u dx &= \int_{\hat{e}} \left[v_1 \cdots v_m \right] \cdot \underbrace{\begin{bmatrix} \phi_1 \\ \vdots \\ \phi_m \end{bmatrix}_{m \times 1}}_{W_e} \underbrace{\left[B_1, B_2 \right] (\Phi) D |J|}_{\underbrace{W_e}_{M_e(\text{mass matrix})}} \begin{bmatrix} u_1 \\ \vdots \\ u_m \end{bmatrix} d\zeta d\eta \\ (III) \int_{\Omega} cvu dx &= \sum_{e \in J_h} \int_e cvu dx \\ \int_{\Omega} cvu dx &= \int_{\hat{e}} \left[v_1 \cdots v_m \right] \underbrace{\begin{bmatrix} \phi_1 \\ \vdots \\ \phi_m \end{bmatrix}}_{M_e(\text{mass matrix})} (\phi_1 \cdots \phi_m) |J| \begin{bmatrix} u_1 \\ \vdots \\ u_m \end{bmatrix} dx \end{aligned}$$

Step 3. Assembling the element matrices K_e, W_e, M_e into global matrices.

$$\sum_{e \in J_h(\Omega)} K_e = K, \quad \sum_{e \in J_h(\Omega)} W_e = W, \quad \sum_{e \in J_h(\Omega)} M_e = M \quad \text{and} \quad \sum_{e \in J_h(\Omega)} \vec{u}_e = \vec{u}$$

here \vec{u}, \vec{v} are the global nodal vector.

The left hand side of (8) is discretized into matrix form

$$\vec{v}(K + W + M)\vec{u} = (I) + (II) + (III) \text{ in (8)}$$

Moreover the righthand side of (8) can also be discretized as

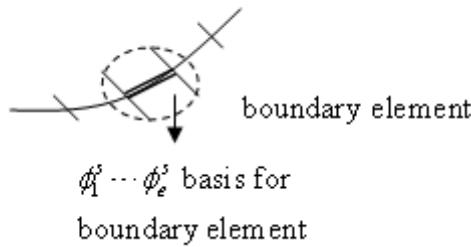
$$(IV)-(8) = \vec{v} \cdot M \cdot \vec{f}$$

$$(V)-(8) = \int_{\Gamma_2} v \cdot r ds = \sum_{e \in J_h(\Gamma_2)} \vec{v}_e \cdot M_e^s \vec{r}_e$$

mass matrix
for boundary
element

$$M_e^s = \int_e \begin{pmatrix} \phi_1^s \\ \vdots \\ \phi_e^s \end{pmatrix} \left(\phi_1^s \cdots \phi_e^s \right) \underbrace{\left| J^s \right|}_{\substack{\text{Jacobian of} \\ \text{surface integral}}} ds = M^s \cdot \vec{r}^s$$

$$\Rightarrow \vec{v} \left(M \vec{f} + M^s \vec{r}^s \right) = (IV) + (V) \text{ in (8)}$$



Step 4. Solve the linear system

$$(K + W + M) \vec{u} = M \vec{f} + M^s \vec{r}^s$$

Example: Given $\Omega = [0, \pi] \times [0, \pi]$, solve

$$\begin{cases} -\Delta u + 2u = 0 \\ u|_{\partial\Omega} = 0 \end{cases} \quad \text{using rectangular 4 nodal element.}$$

where each boundary segment is decomposed into 4, 8, 16, 32 uniform partition.

Compare your answer with the exact solution $u(x, y) = \sin x \sin y$.

